

I | Photoelectric effect 30

①  $W_0 = 4.3 \text{ eV}$   $K_{e^-} = 1 \text{ eV}$

16 •  $E_{ph} = qW_0 + q = \underline{8.5 \cdot 10^{-19} \text{ J}} = \underline{5.3 \text{ eV}}$  4

•  $\nu = \frac{E_{ph}}{h} = \underline{1284 \text{ THz}}$  4

•  $\lambda = \frac{hc}{E_{ph}} = \underline{0.234 \mu\text{m}}$  4

• momentum  $p = \frac{h}{\lambda} = 2.83 \cdot 10^{-27} \frac{\text{kg} \cdot \text{m/s}}{\text{or } 55/\text{me}}$  4

14 ②  $W_0 = 2 \text{ eV}$   
we know that

$K_{e^-} = h(\nu - \nu_0)$  if  $\nu > \nu_0$   
or  $\lambda < \lambda_0$

Threshold  $\nu_0 = \frac{W_0}{h} = \underline{483.5 \text{ THz}}$

or  $\lambda_0 = \frac{c}{\nu_0} = \frac{hc}{W_0} = \underline{620 \text{ nm}}$  4

\* [i] incoming radiation is  $\lambda = 610 \text{ nm}$  (2)

$$\lambda < \lambda_0 \quad \nu > \nu_0$$

$$0.054 \cdot 10^{-19} \text{ J}$$

so  $K_{e^-} = h\nu - h\nu_0 = h\nu - W_0 = \underline{34 \text{ meV}}$  4

~~\*\*\*~~  $K_{e^-} = \frac{p^2}{2m} = \frac{1}{2} m v^2$

$$\Rightarrow \underline{v = \sqrt{\frac{2K_{e^-}}{m}}} = \underline{1.09 \cdot 10^5 \text{ m/s}}$$
 4

\* [i] incoming radiation is  $\lambda = 630 \text{ nm}$  2

$\lambda > \lambda_0$   
 $\nu < \nu_0$  below threshold, photoelectric effect does not happen.

## II Hydrogen atom (30)

$$|\psi(r)| = C \exp\left(-\frac{r}{a_0}\right)$$

①  $\int_0^{+\infty} |\psi|^2 4\pi r^2 dr = 1$  Normalization

4/  $\Rightarrow 4\pi C^2 \int_0^{+\infty} r^2 \exp\left(-\frac{2r}{a_0}\right) dr = 1$

Integration by parts

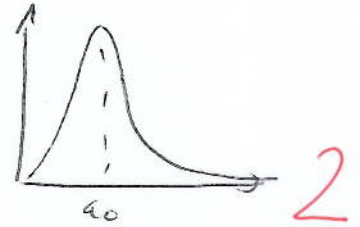
$$\int_0^{+\infty} r^2 \exp\left(-\frac{2r}{a_0}\right) dr = a_0 \int_0^{+\infty} r \exp\left(-\frac{2r}{a_0}\right) dr = \frac{a_0^2}{2} \left[ \exp\left(-\frac{2r}{a_0}\right) \right]_0^{+\infty} = \frac{a_0^3}{4}$$

$$\Rightarrow C = \left(\frac{1}{\pi a_0^3}\right)^{1/2} \quad 4$$

(3)

$$(2) \quad dP = |\psi|^2 4\pi r^2 dr$$

$$\text{so } \frac{dP}{dr} = \left(\frac{4}{a_0^3}\right) \exp\left(-\frac{2r}{a_0}\right) r^2 \quad 4$$



2

(3)  $E(r) = \frac{dP}{dr}$  maximum for  $E(r)$  obtained if.

$$\frac{dE(r)}{dr} = 0 \Rightarrow \underline{r_p = a_0} \quad \text{after calculations. } 4$$

• [The Bohr radius corresponds, in fact, to an orbit where the probability to find the  $e^-$  is maximal. 2

$$(4) \quad P(0.9r_p < r < 1.1r_p) = \left(\frac{4}{a_0^3}\right) \int_{0.9r_p}^{1.1r_p} r^2 \exp\left(-\frac{2r}{a_0}\right) dr$$

After calculations we get  $\underline{P = 0.108} \quad (10.8\%)$  4

⑤  $\langle r \rangle = \frac{4}{a_0^3} \int_0^{+\infty} r^3 \exp\left(-\frac{2r}{a_0}\right) dr = \frac{3a_0}{2}$  after calculation ④

⑥  $U(r) = -\frac{\alpha}{r^2}$  with  $\alpha = \frac{q^2}{4\pi\epsilon_0}$

$\langle U \rangle = -\alpha^2 \langle \frac{1}{r} \rangle = -\alpha^2 \left( \frac{4}{a_0^3} \right) \int_0^{+\infty} r \exp\left(-\frac{2r}{a_0}\right) dr = \frac{\alpha^2}{a_0}$

$\langle U \rangle = -27.1 \text{ eV}$  /2

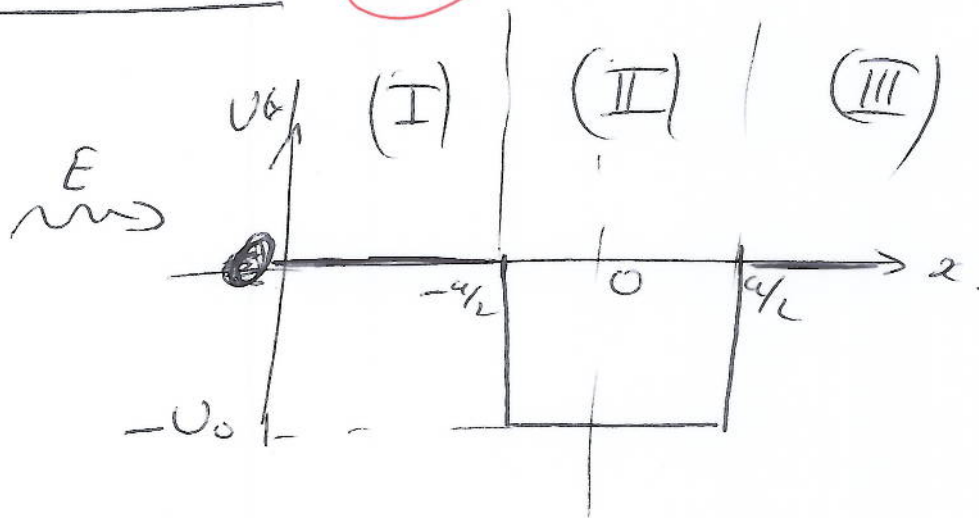
also  $T = \frac{\hbar^2}{2m r^2} \frac{\partial}{\partial r} \left( \frac{\partial \psi}{\partial r} \right)$  (in spherical coordinates)

$\Rightarrow \langle T \rangle = \frac{\hbar^2}{2m} \frac{1}{a_0^2} = 13.6 \text{ eV}$  /2

Since  $E = T + U$ , it comes  $E = -13.6 \text{ eV}$  /2

equivalent to the ground state energy obtained by the Bohr model

# III Finite potential well (25) (5)



①

$$\psi_I = e^{ikx} + r e^{-ikx} \quad /2$$

$$\psi_{II} = E e^{ik_1 x} + F e^{-ik_1 x} \quad /2$$

$$\psi_{III} = t e^{i(k_1 x - \omega t)} \quad /2$$

②  $R + T = 1 \Rightarrow R = 1 - T \quad /2$

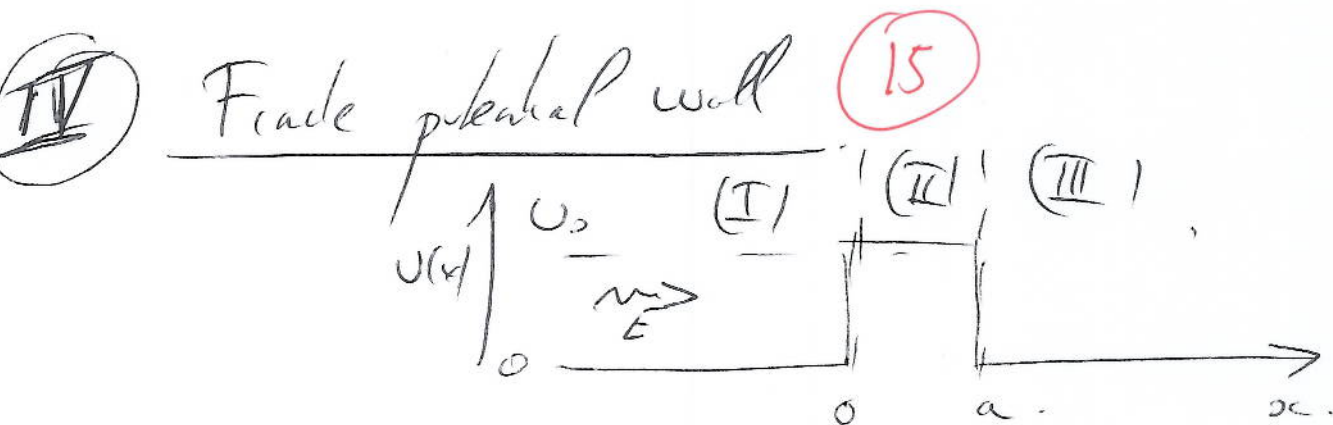
③ see figure for  $T(E) \quad /1$

④  $T = 1$  for  $g(E) = n\pi \quad (n \in \mathbb{N})$ ,

we get  $E_n = \frac{n^2 \pi^2 \hbar^2}{2ma^2} - U_0 \quad /2$

$\Rightarrow E_1 = 0.7 \text{ eV} \quad /1$   
 $E_2 = 50.8 \text{ eV} \quad /1$   
 $E_3 = 134.3 \text{ eV} \quad /1$   
 (see graph ab.)

(5)  $\lim_{E \rightarrow \infty} f(E) \rightarrow 0$   $T \rightarrow 1$  /2 (6)  
 at high energy, the  $e^-$  "does not feel" the effect of the potential well. Same result than that obtained using classical mechanics. /3



(1)  $a \gg \frac{1}{\alpha} \Rightarrow \exp(-\alpha a) \ll \exp(\alpha a)$  /2  
 $\Rightarrow T \approx \frac{16 e^{-2\alpha a}}{\left(4 + \left(\frac{\alpha}{k} - \frac{k}{\alpha}\right)^2\right)}$

$\Rightarrow T \approx \frac{16 E (1-E)}{U_0} e^{-2\alpha a}$  /4

② From previous relation.

$$\text{Log}(T) = \text{Log} \left[ 16 \frac{E}{U_0} \left( 1 - \frac{E}{U_0} \right) \right] - 2\alpha a.$$

⇒ error

$E$

$$\Rightarrow \boxed{\text{Log}(T) \approx -2\alpha a + E}$$

•  $E(x) > 0$  is a necessary condition to obtain max error, indeed  $T < e^{(2\alpha a + E)}$  only if  $e^E > 1 \Rightarrow E > 0$

• we set  $x = \frac{E}{U_0}$ ,  $E(x)$  is maximal if  $\frac{dE}{dx} = 0$ .

$$\bullet \text{ do } \left[ E \left( \frac{1}{4} \right) = \text{Log}(4) \approx 1.4 \right] \Rightarrow x = \frac{1}{2}.$$

when  $\alpha a = 50$  we get  $\text{Log}(T) = -100 + 1.4$

the formula  $\text{Log}(T) = -2\alpha a$  is then exact

at 1.4%

• At the classical limit  $\hbar \rightarrow 0$   $\text{Log} T \rightarrow -\infty$  thus  
 $T \rightarrow 0$

