

# Discrete Calculus Methods and their Application to Fluid Dynamics Blair Perot

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#### Mathematical Tools Order of Accuracy Stability Consistency

Physical Requirements Conservation (mass, momentum, total energy) Secondary Conservation (kinetic energy, vorticity, entropy) Unphysical Modes (pressure) Wave propagation (direction) Eigenmodes/Resonance







# Why do some methods capturephysics well?They use exact discretization

#### What do they have in common? All are Discrete Calculus Methods

Can you make a numerical method so that is mimetic from its design? Yes (I think)



#### Discretization





Take a continuous problem to a finite dimensional one.



#### **Discrete Calculus: Part 1**

**Exact Discretization** 

$$\frac{\partial a}{\partial t} + \nabla \cdot \mathbf{b} = \mathbf{0} \quad \Longrightarrow \quad \left[ \mathbf{A} \quad \mathbf{B} \right] \left( \begin{array}{c} \mathbf{a} \\ \mathbf{b} \end{array} \right] = \vec{\mathbf{r}}$$

Infinite Dimensional Finite Dimensional

Partial Differential Eqn. Matrix Problem

Basic unknowns are integral quantities. Collect infinite data into finite groups.



#### **Discrete Calculus: Part 2**

**Solution** <u>requires</u> Approximation

 $\begin{bmatrix} \mathbf{A} & \mathbf{B} \end{bmatrix} \begin{pmatrix} \vec{\mathbf{a}} \\ \vec{\mathbf{b}} \end{pmatrix} = \vec{\mathbf{r}} \implies \begin{bmatrix} \mathbf{A} & \mathbf{B} \\ \mathbf{C} & \mathbf{D} \end{bmatrix} \begin{pmatrix} \vec{\mathbf{a}} \\ \vec{\mathbf{b}} \end{pmatrix} = \begin{pmatrix} \vec{\mathbf{r}} \\ \mathbf{0} \end{pmatrix}$ 

#### Underdetermined

**Unique Square** 

Relate discrete unknowns to each other. This relation is a material law. Also related to an dual mesh interpolation. Also related to inner products

# Implications

#### Exact Discretization means that:

- Calculus is exact.
- Physics is exact.
- Method is mimetic.

#### **Discrete Solution Requires:**

- Approximation of material laws.
- Interpolation between meshes.
- Assumptions about the solution.



## **Example: Heat Equation**

 $\frac{\partial T}{\partial t} = \alpha \nabla^2 T$ 



Figure out what should NOT be approximated and what is already an approximation (Tonti).



# **Heat Eqn: Fully Separated**

#### **Heat Equation**

 $\frac{\partial i}{\partial t} = -\nabla \cdot \mathbf{q}$ 

 $\mathbf{g} = \nabla T$ 

- Conservation of Energy (Physics)
- $i \approx \rho CT$  Perfectly Caloric Material (Mat.)
- $\mathbf{q} \approx -\mathbf{k}\mathbf{g}$  Fourier's Heat Flux Law (Mat.)
  - Def. of Gradient (Math)
- Discretize Physics and Math exactly.
   Approximate Material laws using interpolation.



Step 1

#### **Exact Discretization**

$I^{n+1}_{ ilde{c}}$ .	$-I_{\tilde{c}}^{n} =$	$-\sum \bar{Q}^{out}_{\tilde{f}}$
		faces

Gauss' Theorem (with time) Like FV: Not very original.

$$I_{\tilde{c}}^{n+1} = \int_{\Omega_{\tilde{c}}^{n+1}} i dV \qquad \overline{Q}_{\tilde{f}}^{out} = \int_{t^n}^{t^{n+1}} dt \int_{\partial \Omega_{\tilde{f}}} \mathbf{q} \cdot \mathbf{n}^{out} dA$$

$$g_{e}^{n+1} = \int_{v_{1}}^{v_{2}} \mathbf{g} \cdot d\mathbf{l} |^{n+1} = \int_{v_{1}}^{v_{2}} \nabla T \cdot d\mathbf{l} |^{n+1} = T_{v_{2}}^{n+1} - T_{v_{1}}^{n+1}$$
  
Math: still exact



#### **Exact Discrete System** $\frac{\partial i}{\partial t}$ $g_e^{n+1}$ $\nabla T$ g $T^{n+1}$

- Exact
- Over-determined
- Uncoupled
- Unknowns on different meshes

Step 2

#### Step 3

#### **Vertex Centered Mimetic**



#### **Approximation Part**

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# $\begin{bmatrix} \mathbf{I} & \mathbf{D} & 0 & 0 \\ \mathbf{I} & 0 & 0 & -\mathbf{M}_{C} \\ 0 & 0 & \mathbf{I} & -\mathbf{G} \\ 0 & \mathbf{I} & \frac{k\Delta t}{2} \mathbf{B}^{n+1} & 0 \end{bmatrix} \begin{bmatrix} I_{\tilde{c}}^{n+1} \\ \overline{Q}_{\tilde{f}} \\ g_{\tilde{f}}^{n+1} \\ T_{v}^{n+1} \end{bmatrix} = \begin{bmatrix} I_{\tilde{v}}^{n} \\ 0 \\ 0 \\ \frac{k\Delta t}{2} \mathbf{B}^{n} g_{e}^{n} \end{bmatrix}$

 $\mathbf{M}_{C}T_{v}^{n+1} - I_{\tilde{c}}^{n} = \mathbf{D}(\frac{\Delta tk}{2}\mathbf{B}^{n+1})\mathbf{G}T_{v}^{n+1} + \mathbf{D}(\frac{\Delta tk}{2}\mathbf{B}^{n})\mathbf{G}T_{v}^{n}$ 

- Single unknown (T at vertices).
- Looks like no dual mesh was used.
  B is Hodge \* operator.

#### **Other Possibilities**

- Mesh: Use polygons as the primary mesh. (quad/hex meshes, particle methods, SOM)
- Approx: Use other basis functions for interpolation. (Rational polynomials, Natural Neighbors, Fourier)
- Exact: Use different exact discretizations (FE, Keller/Priessman Box Schemes)



# **Cell Centered Mimetic**

Q

- Cell Centered
- Median dual mesh
- Raviart-Thomas for Q

 $\overline{I_c^{n+1}} \approx M_c C_c \overline{T_{\tilde{c}}}$  $i \approx \rho \overline{CT}$  $\frac{\Delta t}{2} \left( g_{\tilde{e}}^{n+1} + g_{\tilde{e}}^{n} \right) \approx -\mathbf{A}_{1/k_c}^{n+1} \overline{Q}_f$  $q \approx -kg$ 

# **Cell System**

0 n+1  $I_c^n$ D  $\mathbf{O}$ 0 0  $-\mathbf{M}_{C}$  $\mathbf{O}$ *n*+1 0 G  $\mathbf{O}$  $\mathbf{O}$  $g_{\tilde{e}}^{\pi}$ n+1 $\frac{\Delta t}{2}$  $T^{n+1}$ A

$$\begin{bmatrix} \mathbf{M}_{C} & \mathbf{D} \\ -\mathbf{G} & -\frac{2}{\Delta t} \mathbf{A}_{1/k_{c}}^{n+1} \end{bmatrix} \begin{pmatrix} T_{\tilde{c}}^{n+1} \\ \overline{Q}_{f} \end{pmatrix} = \begin{pmatrix} I_{c}^{n} \\ g_{\tilde{e}}^{n} \end{pmatrix}$$

Symmetric system of unknowns
Not reducible

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#### Comparison



Mimetic 10x more accurate or 1/10<sup>th</sup> cost.

## **Fluid Example - Droplets**



- Moving Mesh Unstructured Surface Tension





# **Example – 3D Droplets**









#### Summary

- Exact discretization (or 2 step construction) leads to mimetic methods.
- Place numerical approximations with the physical approximation.
- Discrete Calculus analysis is accessible to all computational scientist. (Gauss' Theorem).
- Discrete Calculus Methods are not a type of numerical method. The approach can produce <u>some</u> FV, FE, FD, and other methods.

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